Neutron negative charge density:Exclusive-Inclusive Connection

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Outline-

1. How not to and how to analyze electromagnetic form factors- transverse density

- 2. Model independent proton, neutron transverse charge density
- 3. Interpret neutron charge density -relate to high x pdfs

Transverse Charge Densities. Gerald A. Miller, arXiv:1002.0355 [nucl-th] ARNPS

Interpretation of Sachs - $G_E(Q^2)$ is Fourier transform of charge density

 $R^2 = -6\frac{dG_E(Q^2)}{dQ^2}|_{Q^2=0}$

Correct non-relativistic: wave function invariant under Galilean transformation, BUT IT IS WRONG Relativistic : wave function is frame dependent, initial and final states differ interpretation of Sachs FF is wrong Final wave function is **boosted** from initial **Need relativistic treatment**

Toy model GAM, Phys.Rev.C80:045210,2009. Scalar meson M, made of two scalar mesons, m IF (M-2m)/M, small non-relativistic works



Light front, Infinite momentum frame

"Time", $x^+ = x^0 + x^3$, "Evolve", $p^- = p^0 - p^3$ "Space", $x^- = x^0 - x^3$, "Momentum", p^+ (Bjorken) Transverse position, momentum \mathbf{b}, \mathbf{p}

These variables are used in GPDs, TMDs, standard variables transverse boosts in kinematic subgroup

$$\mathbf{k} \to \mathbf{k} - k^+ \mathbf{v}$$

space – like $q^{\mu}, q^{+} = 0$,

momentum transfer in transverse direction

then density is 2 Dimensional Fourier Transform

Model independent transverse charge density

$$J^{+}(x^{-}, \mathbf{b}) = \sum_{q} e_{q} q_{+}^{\dagger}(x^{-}, b) q_{+}(x^{-}, b)$$
Charge Density

$$\rho_{\infty}(x^{-}, \mathbf{b}) = \langle p^{+}, \mathbf{R} = \mathbf{0}, \lambda | \sum_{q} e_{q} q_{+}^{\dagger}(x^{-}, b) q_{+}(x^{-}, b) | p^{+}, \mathbf{R} = \mathbf{0}, \lambda \rangle$$

$$F_{1} = \langle p^{+}, \mathbf{p}', \lambda | J^{+}(0) | p^{+}, \mathbf{p}, \lambda \rangle$$

$$\rho(b) \equiv \int dx^{-} \rho_{\infty}(x^{-}, \mathbf{b}) = \int \frac{Q dQ}{2\pi} F_{1}(Q^{2}) J_{0}(Qb)$$
Density is $u - \bar{u}, \ d - \bar{d}$
Soper '77

Impact parameter dependent GPD Burkardt

Probability that quark at b from CTM has long momentum fraction x

$$\rho^{q}(b,x) = \int \frac{d^{2}q}{(2\pi)^{2}} e^{-i\mathbf{q}\cdot\mathbf{b}} H_{q}(x,t=\mathbf{q}^{2})$$

$$\rho(b) = \sum_{q} e_{q} \int dx \rho^{q}(b,x) \qquad \begin{array}{c} \text{sum rule: integral} \\ \text{of } \mathsf{H}_{q} \quad \text{is F} \end{array}$$

$$\mathbf{R} = \mathbf{0} = \sum_{i}^{N} x_{i} \mathbf{b}_{i}$$

Quark of x=1, must have b=0

Transverse density is integral over longitudinal position or momenta example of Parseval's theorem

Transverse charge densities from parameterizations (Alberico)



Neutron



Neutron charge density: why?



Neutron interpretation

- Impact parameter gpd Burkardt $\,
 ho(x,b)$
- Drell-Yan-West relation between high x DIS and high Q² elastic scattering
- High x related to low b, not uncertainty principle

 $\lim_{x \to 1} \nu W_2(x) = (1-x)^{2n-1} \leftrightarrow \lim_{Q^2 \to \infty} F_1(Q^2) \sim \frac{1}{Q^{2n}}, n = 2$

 Various (reasonable) assumptions needed: power-law wave functions. Is this relationship valid?

$$\pi \text{ DYW}: F(Q^2) \sim \frac{1}{Q^2} \to \nu W_2 = (1-x)$$

But Reimer data $(1-x)^2$???? 10

Structure functions

- d quarks dominate DIS from neutron at high x
- d quarks dominate at neutron center, or π

Density is $u - \bar{u}, d - \bar{d}$ π^- is $\bar{u}d$ decreases u contribution enhances d contribution



Neutron interpretation





Summary

- Much data exist, Jlab12 will improve data set
- Charge density is not a 3 dimensional Fourier transform of G_E
- Interpret form factor as determining transverse charge density
- Neutron: Negative central density
- Full understanding needs $\rho(b, x)$
- Center of neutron: d or π
- Is Drell-Yan, West relation valid?

Spares follow



Relation or not between GPD and TMD

GPD :

$$\langle P', S' | \int \frac{dx^{-}}{4\pi} \bar{q}(-\frac{x^{-}}{2}, \mathbf{0}) \gamma^{+} q(\frac{x^{-}}{2}, \mathbf{0}) e^{ix\bar{p}^{+}x^{-}} | P, S \rangle_{x^{+}} = 0$$

= $\frac{1}{2\bar{p}^{+}} \bar{u}(P', S') \left(\gamma^{+} H_{q}(\xi, t) + i \frac{\sigma^{+\nu} \Delta_{\nu}}{2M} E_{q}(x, \xi, t) \right) u(P, S)$

TMD :

$$\Phi_q^{\Gamma}(x = \frac{k^+}{P^+}, \mathbf{k}) = \langle P, S | \int \frac{d\zeta^- d^2 \zeta}{2(2\pi)^3} e^{ik \cdot \zeta} \bar{q}(0) \Gamma q(\zeta) | P, S \rangle_{\zeta^+ = 0}$$

GPD: nucleons have different momenta, but FT local in coordinate space if integrate over x

TMD: nucleons have same momenta, operator is local in momentum space

Both can be obtained Wigner distribution operator

$$\begin{split} W_q^{\Gamma}(\zeta^-,\boldsymbol{\zeta},k^+,\mathbf{k}) \\ &= \frac{1}{4\pi} \int d\eta^- d^2 \eta e^{i\boldsymbol{k}\cdot\boldsymbol{\eta}} \bar{q}(\zeta^- - \frac{\eta^-}{2},\boldsymbol{\zeta} - \frac{\boldsymbol{\eta}}{2}) \Gamma q(\zeta^- + \frac{\eta^-}{2},\boldsymbol{\zeta} + \frac{\boldsymbol{\eta}}{2}) \\ H_q(x,\xi,t) &= \langle P',S'| \int \frac{d^2\mathbf{k}}{(2\pi)^2} W_q^{\gamma^+}(\zeta^- = 0,\zeta = 0,k^+,\mathbf{k}) | P,S \rangle \\ \Phi_q^{\Gamma}(x,\mathbf{k}) &= \langle P,S| \int \frac{d\zeta^-}{(2\pi)^2} W_q^{\Gamma}(\zeta^-,\boldsymbol{\zeta},k^+,\mathbf{k}) | P,S \rangle \end{split}$$